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A New Split-Weight Scheme for Finite- β Gyrokinetic Plasmas W.W. LEE, E.A. STARTSEV, W.X. WANG, Princeton Plasma Physics Laboratory, Princeton, NJ 08543 — The original split-weight scheme for finite- β simulations [1], which separates the perturbed particle distribution into an adiabatic part and a nonadiabatic part, is generalized to include spatial inhomogeneities. The new scheme requires an additional separation of the fast particle response associated with quasistatic bending of the magnetic field lines. While the original scheme follows the non-adiabatic response, δh , in time, where $\delta h = F - (1 + \psi)F_0$, F is the distribution, F_0 is the background, $\psi \equiv \phi + \int (\partial A_{\parallel}/\partial t) dx_{\parallel}/c$ and ϕ and A_{\parallel} are the perturbed potentials, the new scheme makes use of $\hat{\mathbf{b}} \cdot \nabla(F_0 + \delta g) = 0$, where $\hat{\mathbf{b}} = \hat{\mathbf{b}}_0 + \delta \mathbf{B}/B_0$, and further separates the plasma response as $F = (1 + \psi)F_0 + \delta g + \delta h$, where $\delta g = \int dx_{||} \kappa \cdot (\nabla A_{||} \times \hat{\mathbf{b}}_0)$ and κ is the zeroth order spatial inhomogeneity. The new δh is again followed in time. The results for finite- β stabilization of drift waves and ion temperature gradient modes in slab geometry using the new scheme with a $\beta (\equiv c_s^2/v_A^2)$ as high as 10% and a grid size of the order of the electron skin depth, are in agreement with those discussed in Refs. [2] and [3]. This work is supported by the DoE OASCR Multi-Scale Gyrokinetics (MSG) Project. [1] W. W. Lee, J. Lewandowski, Z. Lin and T. S. Hahm, Phys. Plasmas 8, 4435 (2001). [2] J. V. W. Reynders, Ph. D. Thesis, Princeton University (1992). [3] J. C. Cummings, Ph. D. Thesis, Princeton University (1995).

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